Optimal Poincaré and logarithmic Sobolev constants by decomposition of the energy landscape

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joint work with Georg Menz (Stanford)

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1 Description and Question

2 Main results

3 Sketch of the Proofs

- Local PI and LSI
- Mean-difference estimate



Introduction

Overdamped Langevin dynamics

Hamiltonian $H : \mathbb{R}^n \to \mathbb{R}$ energy landscape

Dynamic at temperature $\varepsilon \ll 1$ d $X_t = -\nabla H(X_t) dt + \sqrt{2\varepsilon} dW_t$

Fokker-Planck evolution of law $X_t = \varrho_t$ $\partial_t \varrho_t = \nabla \cdot (\varepsilon \nabla \varrho_t + \varrho_t \nabla H)$

$$\begin{array}{ll} \text{Gibbs measure } \mu(\mathsf{d} x) = \frac{1}{Z_{\mu}}\exp\left(-\frac{H}{\varepsilon}\right)\mathsf{d} x,\\ \text{where } \quad Z_{\mu} = \int \exp\left(-\frac{H}{\varepsilon}\right)\,\mathsf{d} x \end{array}$$

Generator evolution of $f_t = \rho_t / \mu$ $\partial_t f_t = L f_t := \varepsilon \Delta f_t - \nabla H \cdot \nabla f_t$







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 μ satisfies the Poincaré inequality $\mathsf{Pl}(\varrho)$ if $\forall f : \mathbb{R}^n \to \mathbb{R}$

$$\mathsf{var}_{\mu}(f) := \int f^2 - \left(\int f \mathsf{d}\mu\right)^2 \mathsf{d}\mu \leq \frac{1}{\varrho} \int |
abla f|^2 \,\mathsf{d}\mu. \qquad \mathsf{PI}(\varrho)$$

and the logarithmic Sobolev inequality $\mathsf{LSI}(\alpha)$ if $\forall f : \mathbb{R}^n \to \mathbb{R}$

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 $PI(\rho)$ and $LSI(\alpha)$ imply exponential convergence to μ :

$$\begin{aligned} \mathsf{PI}(\varrho) \ \Rightarrow \ \mathsf{var}_{\mu}(f_t) \leq \mathsf{var}_{\mu}(f_0) e^{-2\varrho\varepsilon t} \\ \mathsf{LSI}(\alpha) \ \Rightarrow \ \mathsf{Ent}_{\mu}(f_t) \leq \mathsf{Ent}_{\mu}(f_0) e^{-2\alpha\varepsilon t} \end{aligned}$$



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Goal: Optimal constants in PI and LSI



Accurate estimates of ϱ and α in the regime $\varepsilon \ll 1$:

$$arrho = \mathcal{C}_{arrho}(arepsilon) e^{-rac{\Delta H}{arepsilon}}(1+o(1))$$
 and $lpha = \mathcal{C}_{lpha}(arepsilon) e^{-rac{\Delta H}{arepsilon}}(1+o(1)).$

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$$\varrho = C_{\varrho}(\varepsilon)e^{-\frac{\Delta H}{\varepsilon}}(1+o(1))$$
 and $\alpha = C_{\alpha}(\varepsilon)e^{-\frac{\Delta H}{\varepsilon}}(1+o(1)).$



$$\mathrm{d}X_t = -
abla H(X_t) \,\mathrm{d}t + \sqrt{2arepsilon} \,\mathrm{d}W_t$$

- particle follows $-\nabla H$ as long as $|\nabla H| \sim 1$
- noise is dominant, if $|\nabla H| \lesssim \sqrt{\varepsilon}$





Figure : Trajectory for $\varepsilon = 0.4$

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Figure : Trajectory for $\varepsilon = 0.2$





Figure : Trajectory for $\varepsilon = 0.1$

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Figure : Trajectory for $\varepsilon = 0.05$ (red $\varepsilon = 0$)

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Two scales by decomposition à la $[GOVW09]^1$



- The partition $\biguplus_i \Omega_i = \mathbb{R}^n$ is called admissible for μ if:
 - (i) For each local minimum $m_i \in \mathcal{M}$ exists $\Omega_i \in \mathcal{P}_{\mathcal{M}}$ with $m_i \in \Omega_i$
 - (ii) The partition sum of each Ω_i is approximately Gaussian

$$u(\Omega_i)Z_{\mu} = rac{(2\pi\varepsilon)^{rac{n}{2}}}{\sqrt{\det
abla^2 H(m_i)}} \exp\left(-rac{H(m_i)}{arepsilon}
ight) (1+o(1))\,.$$

Restricted measures: $\mu_i := \mu \llcorner \Omega_i$, i = 0, 1.

Macroscopic measures $\bar{\mu}$ on $\{0, 1\}$: $\bar{\mu} := Z_0 \delta_0 + Z_1 \delta_1.$

Mixture representation:

$$\mu = Z_0\mu_0 + Z_1\mu_1$$
 with $Z_i := \mu(\Omega_i)$.



¹N. Grunewald, F. Otto, C. Villani, and M. G. Westdickenberg, A^s twö-scäle appröach^ato ¹⁵ logarithmic Sobolev inequalities and the hydrodynamic limit, Annales de l'Institut Henri Poincaré, Probabilités et Statistiques, 45:**2**, 2009.

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Splitting

Ideas motivated from [CM10]²

$$\operatorname{var}_{\mu}(f) = \underbrace{Z_{0} \operatorname{var}_{\mu_{0}}(f) + Z_{1} \operatorname{var}_{\mu_{1}}(f)}_{\operatorname{local variances}} + \underbrace{Z_{0} Z_{1} \underbrace{\left(\mathbb{E}_{\mu_{0}}(f) - \mathbb{E}_{\mu_{1}}(f)\right)^{2}}_{\operatorname{mean-difference}}}_{\operatorname{mean-difference}}$$

$$\operatorname{Ent}_{\mu}(f^{2}) = \underbrace{Z_{0} \operatorname{Ent}_{\mu_{0}}(f^{2}) + Z_{1} \operatorname{Ent}_{\mu_{1}}(f^{2})}_{\operatorname{Ent}_{\mu_{1}}(f^{2})} + \underbrace{\operatorname{Ent}_{\mu}\left(\mathbb{E}_{\mu_{\bullet}}(f^{2})\right)}_{\operatorname{Ent}_{\mu}\left(\mathbb{E}_{\mu_{\bullet}}(f^{2})\right)}$$

where
$$\Lambda(Z_0, Z_1) = \frac{Z_0 - Z_1}{\log Z_0 - \log Z_1}$$
 is the logarithmic mean.

Expect from heuristics:

- good estimate for local variances/entropies
- exponential estimate for mean-difference

²D. Chafaï and F. Malrieu, *On fine properties of mixtures with respect to concentration of measure and Sobolev type inequalities*, Annales de l'Institut Henri Poincaré, Probabilités et Statistiques, 46:**1**, 2010.

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- Local PI and LSI
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Theorem (Local PI and LSI)

There exists an admissible partition $\biguplus_i \Omega_i = \mathbb{R}^n$ such that each local measure $\mu_i = \mu_{\perp}\Omega_i$ satisfies $\mathsf{PI}(\varrho_{\mathsf{loc}})$ and $\mathsf{LSI}(\alpha_{\mathsf{loc}})$ with

$$\varrho_{loc}^{-1} = O(\varepsilon) \quad \text{and} \quad \alpha_{loc}^{-1} = O(1).$$

PI is as good as for convex potential

Non-convexity of potential worsens LSI

• Both results scale optimal in one dimension

Theorem (Mean-difference estimate)

 $(\mathbb{E}_{\mu_0}f - \mathbb{E}_{\mu_1}f)^2 \lesssim \frac{Z_{\mu}}{(2\pi\varepsilon)^{\frac{\ell}{2}}} \; \frac{2\pi\varepsilon\sqrt{|\det\nabla^2 H(s_{0,1})|}}{|\lambda^-(\nabla^2 H(s_{0,1}))|} \; e^{\varepsilon^{-1}H(s_{0,1})} \int |\nabla f|^2 \, \mathrm{d}\mu.$

 \lesssim ": up to multiplicative error 1+o(1) as arepsilon o 0.



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" \lesssim ": up to multiplicative error 1 + o(1) as $\varepsilon \to 0$.

Eyring-Kramers formula



Corollary

The measure μ satisfies $PI(\varrho)$ and $LSI(\alpha)$ with

$$\frac{1}{\varrho} \approx Z_0 Z_1 \frac{Z_\mu}{(2\pi\varepsilon)^{\frac{n}{2}}} \ \frac{2\pi\varepsilon\sqrt{|\det\nabla^2 H(s_{0,1})|}}{|\lambda^-(\nabla^2 H(s_{0,1}))|} e^{\frac{H(s_{0,1})}{\varepsilon}} \quad \text{and} \quad \frac{2}{\alpha} \approx \frac{1}{\Lambda(Z_0,Z_1) \ \varrho}.$$

Asymptotic evaluation of the factor $\Lambda(Z_0,Z_1)$ for two special cases:

$$H(m_0) < H(m_1): \quad \frac{\varrho}{\alpha} \approx \frac{1}{2} \left(\frac{H(m_1) - H(m_0)}{\varepsilon} + \log\left(\frac{\kappa_0}{\kappa_1}\right) \right) = O(\varepsilon^{-1})$$
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 ⇒ rules out Bakry-Émery criterion
- non-exponential behavior of constants
 ⇒ rules out Holley-Stroock perturbation principle
- optimality available in one dimension
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Proof: Local PI and LSI via Lyapunov condition universitation important in the second second

Technique developed by Bakry, Barthe, Cattiaux, Guillin, Wang and Wu 2008-

Definition (Lyapunov condition on domains)

L satisfies a Lyapunov condition with constants $\lambda, b > 0$ and some $U \subset \Omega$, if there exists a Lyapunov function $W : \Omega \to [1, \infty)$ satisfying

$$\frac{LW}{\varepsilon W} \leq -\lambda + b \, \mathbb{1}_U.$$

and Neumann boundary condition on Ω , such that integration by parts holds $\int_{\Omega} f(-LW) \, \mathrm{d}\mu = \varepsilon \int_{\Omega} \langle \nabla f, \nabla W \rangle \, \mathrm{d}\mu.$

Theorem ([BBCG08])

Suppose L satisfies a Lyapunov condition and $\mu \sqcup U$ satisfies $PI(\varrho_U)$, then μ satisfies $PI(\varrho)$ with λ

$$\varrho \ge \frac{\lambda}{b + \varrho \upsilon} \varrho \upsilon$$

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Technique developed by Bakry, Barthe, Cattiaux, Guillin, Wang and Wu 2008-

Definition (Lyapunov condition on domains)

L satisfies a Lyapunov condition with constants $\lambda, b > 0$ and some $U \subset \Omega$, if there exists a Lyapunov function $W : \Omega \to [1, \infty)$ satisfying

$$\frac{LW}{\varepsilon W} \leq -\lambda + b \, \mathbb{1}_U.$$

and Neumann boundary condition on Ω , such that integration by parts holds $\int_{\Omega} f(-LW) \, \mathrm{d}\mu = \varepsilon \int_{\Omega} \langle \nabla f, \nabla W \rangle \, \mathrm{d}\mu.$

Theorem ([BBCG08])

Suppose L satisfies a Lyapunov condition and $\mu \sqcup U$ satisfies $PI(\varrho_U)$, then μ satisfies $PI(\varrho)$ with λ

$$\varrho \geq \frac{\lambda}{b + \varrho \upsilon} \varrho \upsilon$$



Proof: Lyapunov \Rightarrow PI(ϱ)

Integration by parts of W wrt. to L yields

$$\begin{split} \int_{\Omega} f^2 \frac{(-LW)}{\varepsilon W} \mathrm{d}\mu &= \int_{\Omega} \left\langle \nabla \frac{f^2}{W}, \nabla W \right\rangle \mathrm{d}\mu \\ &= 2 \int_{\Omega} \frac{f}{W} \left\langle \nabla f, \nabla W \right\rangle \mathrm{d}\mu - \int_{\Omega} \frac{f^2 \left| \nabla W \right|^2}{W^2} \mathrm{d}\mu \\ &= \int_{\Omega} \left| \nabla f \right|^2 \mathrm{d}\mu - \int_{\Omega} \left| \nabla f - \frac{f}{W} \nabla W \right|^2 \mathrm{d}\mu. \end{split}$$

The Lyapunov conditions ensures $1 \leq \frac{-LW}{\lambda \in W} + \frac{b}{\lambda} \mathbb{1}_U$:

$$\operatorname{var}_{\mu}(f) = \int_{\Omega} (f - \overline{f})^2 \mathrm{d}\mu$$



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Lyapunov function

• Task: Find a function $W: \Omega \to [1,\infty)$ such that

$$rac{LW}{W} \leq -\lambda + b \ \mathbb{1}_{B_{a\sqrt{arepsilon}}(m)}.$$

• Ansatz $W = \exp\left(\frac{\tilde{H}}{2\varepsilon}\right)$, where \tilde{H} is an ε -perturbation of H $\frac{\tilde{L}W}{W} = \frac{1}{2}\Delta \tilde{H} - \frac{1}{4\varepsilon}|\nabla \tilde{H}|^2 \stackrel{!}{\leq} -\lambda.$

if x is √ε-away from critical points: ε⁻¹|∇H̃(x)|² ≥ 4λ
 if x is √ε-nearby a critical point of index k ≥ 1

$$\Delta \tilde{H}(x) = \underbrace{\tilde{\lambda}_1^- + \dots + \tilde{\lambda}_k^-}_{<0} + \underbrace{\tilde{\lambda}_{k+1}^+ + \dots + \tilde{\lambda}_n^+}_{>0} + O(\sqrt{\varepsilon})$$



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Construction of Lyapunov function



Figure : *H* around a saddle point

 $ilde{H}$ is quadratic perturbation of H in $\sqrt{arepsilon}$ -neighborhoods of critical points:

$$\sup_{x} \left| H(x) - \tilde{H}(x) \right| = O(\varepsilon).$$



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Figure : *H* around a saddle point

Figure : \tilde{H} around a saddle point

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Approximation step

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Step 1: Approximate μ_0 and μ_1 by truncated Gaussians ν_0 and ν_1 :

 $u_i \sim \mathcal{N}(m_i, \varepsilon \Sigma_i) \llcorner B_{\sqrt{\varepsilon}}(m_i) \text{ with } \Sigma_i^{-1} := \nabla^2 H(m_i).$

Introduce ν_0 and ν_1 as coupling:

$$\left(\mathbb{E}_{\mu_0}f - \mathbb{E}_{\mu_1}f\right)^2 \le (1+\tau)\underbrace{\left(\mathbb{E}_{\nu_0}f - \mathbb{E}_{\nu_1}f\right)^2}_{\bullet}$$

transport argument

$$+2(1+ au^{-1})\sum_{i=\{0,1\}} \underbrace{(\mathbb{E}_{\mu_i}f-\mathbb{E}_{
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Approximation bound follows from local PI and local LSI.

André Schlichting (IAM Bonn)

Optimal PI and LSI



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Sideremark: Weighted transport distance

Definition

For $\nu_0, \nu_1 \ll \mu$ define the weighted transport distance by

$$\mathcal{T}^2_{\mu}(\nu_0,\nu_1) = \inf_{\{\Phi_s\}} \int \left| \int_0^1 \dot{\Phi}_s \circ \Phi_s^{-1} \; \frac{\mathsf{d}\nu_s}{\mathsf{d}\mu} \; \mathsf{d}s \right|^2 \mathsf{d}\mu.$$

 $(\Phi_s)_{s\in[0,1]}$ is absolutely continuous in s: $(\Phi_s)_{\sharp}\nu_0 = \nu_s$.

Mean-difference revisited: Identify $\int |
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$$\left(\int f \, \mathrm{d}\nu_0 - \int f \, \mathrm{d}\nu_1\right)^2 = \left(_{\dot{H}^{-1}(\mu)} \langle \nu_0 - \nu_1, f \rangle_{\dot{H}^1(\mu)}\right)^2 \\ \leq \mathcal{T}^2_{\mu}(\nu_0, \nu_1) \, \|f\|^2_{\dot{H}^1(\mu)} \,.$$

Indeed, it holds: $\mathcal{T}^2_{\mu}(
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Construction of transport interpolation

Step 3: Ansatz
$$\Phi_s$$
 such that $\nu_s = (\Phi_s)_{\sharp} \nu_0 = \mathcal{N}(\gamma_s, \Sigma_s) \sqcup B_{\sqrt{\varepsilon}}(\gamma_s)$

(1) optimize $\gamma \Rightarrow$ passage of saddle $\gamma_{\tau^*} = s_{0,1}$ (2) optimize $\dot{\gamma}_{\tau^*} \Rightarrow$ direction of eigenvector to $\lambda^-(\nabla^2 H(s_{0,1}))$ (3) optimize $\Sigma_{\tau^*} \Rightarrow \Sigma_{\tau^*}^{-1} = \nabla^2 H(s_{0,1})$ on stable manifold of $s_{0,1}$





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optimize γ ⇒ passage of saddle γ_{τ*} = s_{0,1}
 optimize γ_{τ*} ⇒ direction of eigenvector to λ⁻(∇²H(s_{0,1}))
 optimize Σ_{τ*} ⇒ Σ_τ⁻¹ = ∇²H(s_{0,1}) on stable manifold of s_{0,1}





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2 Main results

3 Sketch of the Proofs

- Local PI and LSI
- Mean-difference estimate





Investigation of specific energy landscape:



Features:

- two global minima *a*, *b*
- additional local minimum c
- saddle points $s_{a,b}$, $s_{a,c}$, $s_{c,b}$

degenerated:

$$|H(s_{a,b}) - H(s_{a,c})| = \delta$$
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Fluxes and reaction rates

Interpretation as chemical reactions

•
$$A = \{|x - a| \le r\}$$
 reactant

•
$$B = \{|x - a| \le r\}$$
 product

• $C = \{|x - a| \le r\}$ intermediate product

What are typical reaction rates and paths?



How to define the reaction rate?

Steady state with inflow of reactants and outflow of products:

 $Lh_{A,B} = 0$ in $(A \cup B)^c$ and $h_{A,B} = \mathbb{1}_A$ in $A \cup B$.

Definition (Reaction rate)

$$k_{A,B} := \varepsilon \int |J_{A,B}|^2 d\mu = \varepsilon \int |\nabla h_{A,B}|^2 d\mu.$$

André Schlichting (IAM Bonn)

Optimal PI and LS



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Entropic switching



Numerics: P. Metzner, C. Schütte and E. Vanden-Eijnden 2006 [MSVE06]



[MSVE06] P. Metzner, C. Schütte, and E. Vanden-Eijnden, *Illustration of transition path theory on a collection of simple examples.* The Journal of chemical physics, 125:**8**, 2006.



series and parallel law



• minima become nodes

- saddles become resistors
- boundary condition becomes voltage source

series and parallel law

$$\frac{1}{R} = \frac{1}{R_{AB}} + \frac{1}{R_{AC} + R_{CB}}$$



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Result and comparison with numerical data



Identification can be justified using the weighted transport distance

Series and parallel law for transport cost

$$k_{A,B} pprox rac{1}{\mathcal{T}_{\mu}^2(\mu\llcorner A, \mu\llcorner B)} pprox rac{1}{\mathcal{T}_{AB}} + rac{1}{\mathcal{T}_{AC} + \mathcal{T}_{CB}},$$

where

$$T_{AB} = \inf_{\Phi \in \Pi(s_{AB})} \int \left| \int_0^1 \dot{\Phi}_s \circ \Phi_s^{-1} \left| \frac{\mathrm{d}\nu_s}{\mathrm{d}\mu} \right|^2 \mathrm{d}\mu.$$

 $\Pi(s_{AB})$: transport interpolations between $\mu \llcorner A$ and $\mu \llcorner B$ across s_{AB} .

| | $\varepsilon = 0.15$ | |
|--------------------------------|----------------------|-----------------------|
| [MSVE06] TPT, flux | $9.47 	imes 10^{-8}$ | $1.912 	imes 10^{-2}$ |
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- Partitions and splitting induced from dynamic (two scales)
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 - good local mixing
 Lyapunov technique handles non-convex situations
 - sharp estimates of mean-difference
 transport representation of A⁻¹-norm and optimization
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