

Eyring-Kramers formula
for
Poincaré
and
logarithmic Sobolev inequalities

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joint work with Georg Menz (Stanford)

Bielefeld Stochastics Afternoon

November 14, 2012



Overdamped Langevin dynamics

Hamiltonian $H : \mathbb{R}^n \rightarrow \mathbb{R}$ *energy landscape*

Dynamic at temperature $\varepsilon \ll 1$

$$dX_t = -\nabla H(X_t)dt + \sqrt{2\varepsilon} dW_t$$

Fokker-Planck evolution of law $X_t = \varrho_t$

$$\partial_t \varrho_t = \nabla \cdot (\varepsilon \nabla \varrho_t + \varrho_t \nabla H)$$

Gibbs measure $\mu(dx) = \frac{1}{Z_\mu} \exp\left(-\frac{H}{\varepsilon}\right) dx$,

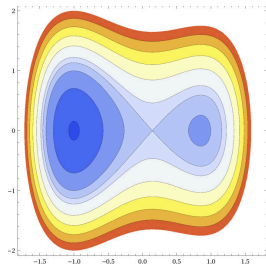
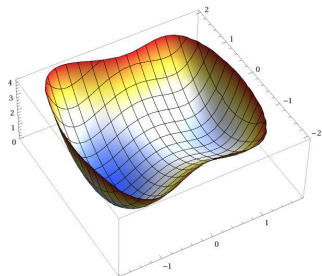
$$\text{where } Z_\mu = \int e^{-\frac{H}{\varepsilon}} dx$$

Generator evolution of $f_t = \varrho_t / \mu$

$$\partial_t f_t = Lf_t := \varepsilon \Delta f_t - \nabla H \cdot \nabla f_t.$$

Dirichlet form $\mathcal{E}(f) := \int (-Lf)f d\mu$

$$= \varepsilon \int |\nabla f|^2 d\mu.$$



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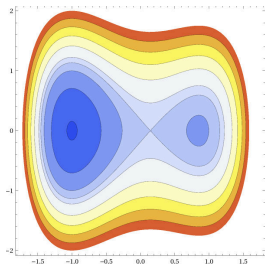
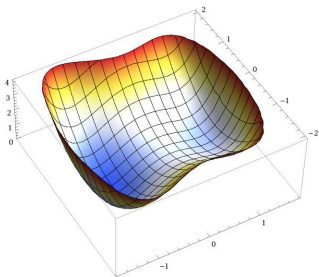
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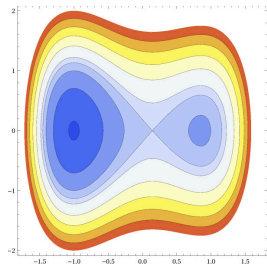
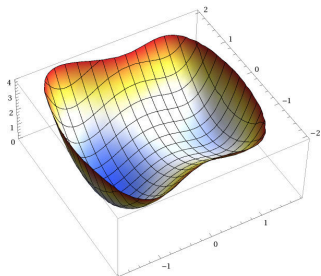
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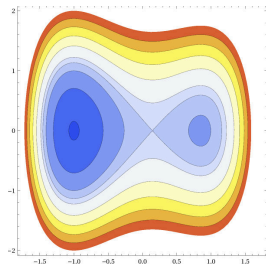
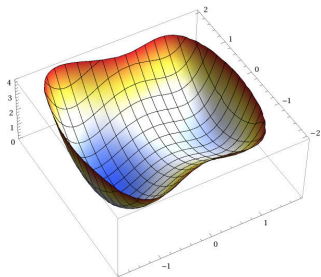
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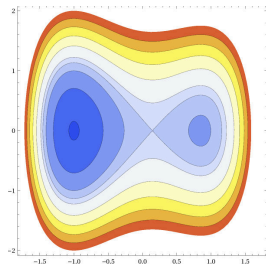
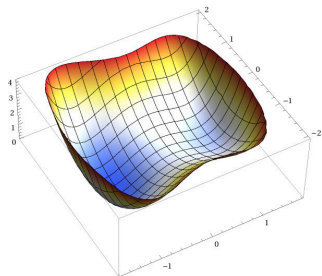
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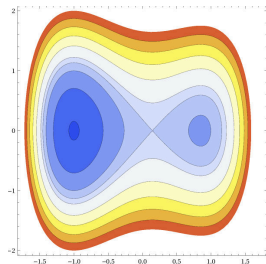
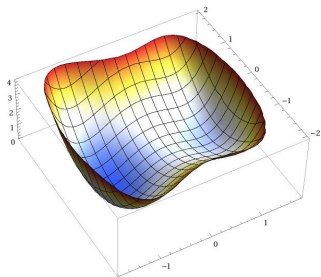
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Quantification via functional inequalities

- $f \equiv \text{const.}$ is equilibrium state
- for a **strictly convex** function $\xi : \mathbb{R} \rightarrow \mathbb{R} \cup \{+\infty\}$ define

$$\Xi(f) := \int \xi \circ f \, d\mu - \xi \left(\int f \, d\mu \right).$$

- evaluate $\Xi(f_t)$ along solution $\partial_t f_t = Lf_t$

$$\frac{d}{dt} \Xi(f_t) = \int \xi' \circ f \underbrace{\partial_t f_t}_{=Lf_t} \, d\mu = -\varepsilon \int \underbrace{\xi'' \circ f}_{>0} |\nabla f_t|^2 \, d\mu.$$

- If **FI**(β), then

$$\Xi(f_t) \leq \Xi(f_0) e^{-2\varepsilon\beta t}.$$

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*Convergence to equilibrium is established by **FI**(β)*

Definition

μ satisfies the **Poincaré inequality** $PI(\varrho)$ if $\forall f : \mathbb{R}^n \rightarrow \mathbb{R}$

$$\text{var}_\mu(f) := \int f^2 - \left(\int f d\mu \right)^2 d\mu \leq \frac{1}{\varrho} \int |\nabla f|^2 d\mu. \quad PI(\varrho)$$

and the **logarithmic Sobolev inequality** $LSI(\alpha)$ if $\forall f : \mathbb{R}^n \rightarrow \mathbb{R}$

$$\text{Ent}_\mu(f) := \int f \log \frac{f}{\int f d\mu} d\mu \leq \frac{1}{\alpha} \int \frac{|\nabla f|^2}{2f} d\mu. \quad LSI(\alpha)$$

$PI(\varrho)$ and $LSI(\alpha)$ imply exponential convergence to μ :

$$PI(\varrho) \Rightarrow \text{var}_\mu(f_t) \leq \text{var}_\mu(f_0) e^{-2\varrho \varepsilon t}$$

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Accurate estimates of ϱ and α in the regime $\varepsilon \ll 1$:

$$\varrho = C_{\varrho}(\varepsilon)e^{-\frac{\Delta H}{\varepsilon}}(1 + o(1)) \quad \text{and} \quad \alpha = C_{\alpha}(\varepsilon)e^{-\frac{\Delta H}{\varepsilon}}(1 + o(1)).$$

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- particle follows $-\nabla H$ as long as $|\nabla H| \sim 1$
- noise is dominant, if $|\nabla H| \lesssim \sqrt{\varepsilon}$

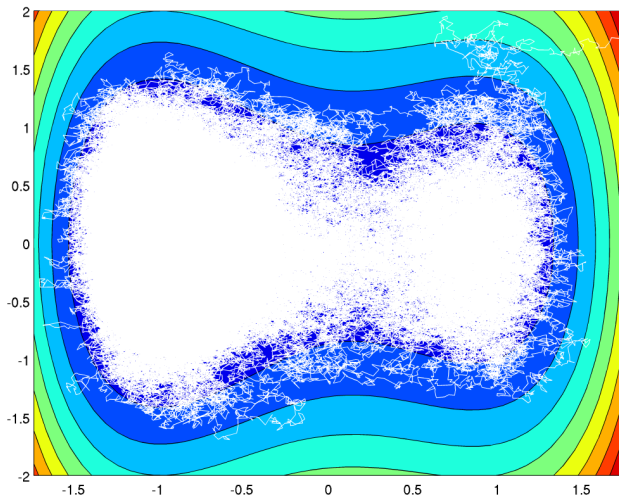


Figure : Trajectory for $\varepsilon = 0.4$

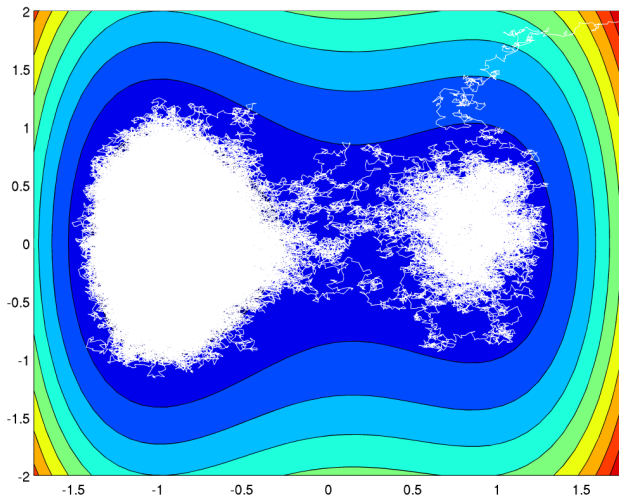


Figure : Trajectory for $\varepsilon = 0.2$

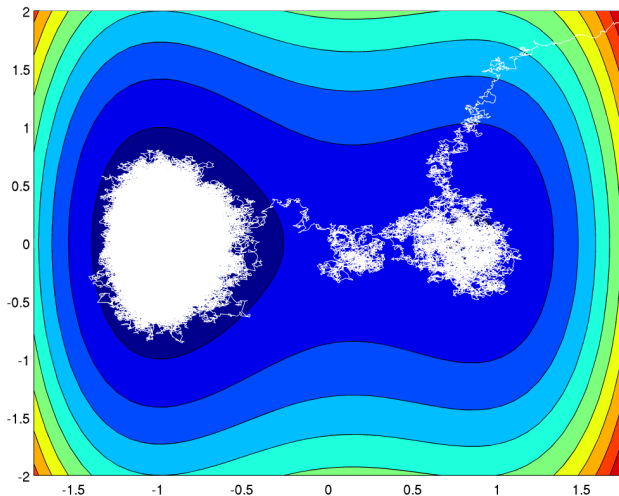


Figure : Trajectory for $\varepsilon = 0.1$

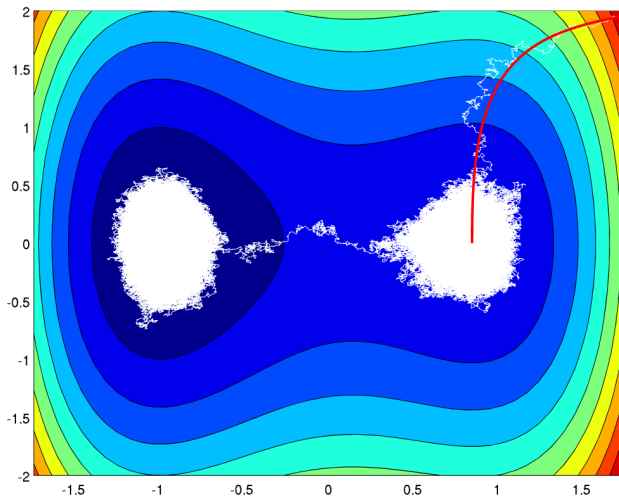


Figure : Trajectory for $\varepsilon = 0.05$ (red $\varepsilon = 0$)

Make use of the two scales by decomposition [GOVW09]¹

Basins of attraction $\Omega_0 \uplus \Omega_1 = \mathbb{R}^n$ of local minima m_0, m_1 :

$$\Omega_i := \{y_0 \in \mathbb{R}^n : \dot{y}_t = -\nabla H(y_t), y_t \rightarrow m_i\}.$$

Restricted measures μ_0, μ_1 :

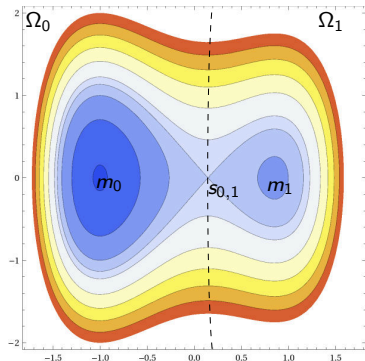
$$\mu_i := \mu \llcorner \Omega_i, \quad i = 0, 1.$$

Macroscopic measures $\bar{\mu}$ on $\{0, 1\}$:

$$\bar{\mu} := Z_0 \delta_0 + Z_1 \delta_1.$$

Mixture representation

$$\mu = Z_0 \mu_0 + Z_1 \mu_1, \quad Z_i := \mu(\Omega_i).$$



¹N. Grunewald, F. Otto, C. Villani, and M. G. Westdickenberg, *A two-scale approach to logarithmic Sobolev inequalities and the hydrodynamic limit*, Annales de l'Institut Henri Poincaré, Probabilités et Statistiques, 45:2, 2009.

Lemma

$$\text{var}_\mu(f) = \underbrace{Z_0 \text{var}_{\mu_0}(f) + Z_1 \text{var}_{\mu_1}(f)}_{\substack{\text{local variances} \\ \text{local entropies}}} + Z_0 Z_1 \underbrace{(\mathbb{E}_{\mu_0}(f) - \mathbb{E}_{\mu_1}(f))^2}_{\text{mean-difference}}$$

$$\begin{aligned} \text{Ent}_\mu(f^2) &\leq \underbrace{Z_0 \text{Ent}_{\mu_0}(f^2) + Z_1 \text{Ent}_{\mu_1}(f^2)} \\ &\quad + \frac{Z_0 Z_1}{\Lambda(Z_0, Z_1)} \left(\text{var}_{\mu_0}(f) + \text{var}_{\mu_1}(f) + (\mathbb{E}_{\mu_0}(f) - \mathbb{E}_{\mu_1}(f))^2 \right), \end{aligned}$$

where $\Lambda(Z_0, Z_1) = \frac{Z_0 - Z_1}{\log Z_0 - \log Z_1}$ is the *logarithmic mean*.

Expect from heuristics:

- *good* estimate for local variances/entropies
- *exponential* estimate for mean-difference

Similar ideas used in [CM10]: D. Chafaï and F. Malrieu, *On fine properties of mixtures with respect to concentration of measure and Sobolev type inequalities*, Annales de l'Institut Henri Poincaré, Probabilités et Statistiques, 46:1, 2010.

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$$\begin{aligned} \text{Ent}_\mu(f^2) &\leq \underbrace{Z_0 \text{Ent}_{\mu_0}(f^2) + Z_1 \text{Ent}_{\mu_1}(f^2)} \\ &\quad + \frac{Z_0 Z_1}{\Lambda(Z_0, Z_1)} \left(\text{var}_{\mu_0}(f) + \text{var}_{\mu_1}(f) + (\mathbb{E}_{\mu_0}(f) - \mathbb{E}_{\mu_1}(f))^2 \right), \end{aligned}$$

where $\Lambda(Z_0, Z_1) = \frac{Z_0 - Z_1}{\log Z_0 - \log Z_1}$ is the *logarithmic mean*.

Expect from heuristics:

- *good* estimate for local variances/entropies
- *exponential* estimate for mean-difference

Similar ideas used in [CM10]: D. Chafaï and F. Malrieu, *On fine properties of mixtures with respect to concentration of measure and Sobolev type inequalities*, Annales de l'Institut Henri Poincaré, Probabilités et Statistiques, 46:1, 2010.

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The measures μ_0 and μ_1 satisfy PI(ϱ_{loc}) and LSI(α_{loc}) with

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$$(\mathbb{E}_{\mu_0} f - \mathbb{E}_{\mu_1} f)^2 \lesssim \frac{Z_\mu}{(2\pi\varepsilon)^{\frac{d}{2}}} \frac{2\pi\varepsilon \sqrt{|\det \nabla^2 H(s_{0,1})|}}{|\lambda^-(\nabla^2 H(s_{0,1}))|} e^{\varepsilon^{-1} H(s_{0,1})} \int |\nabla f|^2 d\mu.$$

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New proof to [BEGK04/05] for PI and extension to LSI:

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The measure μ satisfies PI(ϱ) and LSI(α) with

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Asymptotic evaluation of the factor $\Lambda(Z_0, Z_1)$ for two special cases:

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Lyapunov condition

Technique developed by Bakry, Barthe, Cattiaux, Guillin, Wang and Wu 2008–
Principal eigenvalue characterization for L by Donsker-Varadhan 1975

Definition

L satisfies a **Lyapunov condition** with constants $\lambda, b > 0$ and some $U \subset \mathbb{R}^n$, if there exists a function $W : \Omega \rightarrow [1, \infty)$ satisfying

$$\frac{LW}{\varepsilon W} \leq -\lambda + b \mathbb{1}_U.$$

W is called **Lyapunov function** for L .

Theorem ([BBCG08])

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Proof: Lyapunov \Rightarrow PI(ϱ)

With the symmetry of $\varepsilon^{-1}(-L)$ in $L^2(\mu)$ follows

$$\begin{aligned}\int f^2 \frac{(-LW)}{\varepsilon W} d\mu &= \int \left\langle \nabla \frac{f^2}{W}, \nabla W \right\rangle d\mu \\ &= 2 \int \frac{f}{W} \langle \nabla f, \nabla W \rangle d\mu - \int \frac{f^2 |\nabla W|^2}{W^2} d\mu \\ &= \int |\nabla f|^2 d\mu - \int \left| \nabla f - \frac{f}{W} \nabla W \right|^2 d\mu.\end{aligned}$$

The Lyapunov conditions ensures $1 \leq \frac{-LW}{\lambda \varepsilon W} + \frac{b}{\lambda} \mathbb{1}_U$:

$$\text{var}_\mu(f) = \int (f - \bar{f})^2 d\mu$$

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Lyapunov function

- **Task:** Find a function $W : \Omega \rightarrow [1, \infty)$ such that

$$\frac{LW}{W} \leq -\lambda + b \mathbb{1}_{B_{a\sqrt{\varepsilon}}(m)}.$$

- Ansatz $W = \exp\left(\frac{\tilde{H}}{2\varepsilon}\right)$, where \tilde{H} is an ε -perturbation of H

$$\frac{\tilde{L}W}{W} = \frac{1}{2}\Delta\tilde{H} - \frac{1}{4\varepsilon}|\nabla\tilde{H}|^2 \stackrel{!}{\leq} -\lambda.$$

- ▶ if x is $\sqrt{\varepsilon}$ -away from critical points: $\varepsilon^{-1}|\nabla\tilde{H}(x)|^2 \geq 4\lambda$
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$$\Delta\tilde{H}(x) = \underbrace{\tilde{\lambda}_1^- + \dots + \tilde{\lambda}_k^-}_{<0} + \underbrace{\tilde{\lambda}_{k+1}^+ + \dots + \tilde{\lambda}_n^+}_{>0} + O(\sqrt{\varepsilon})$$

Can negative eigenvalues be enforced such that $\Delta\tilde{H}(x) \leq -2\lambda$?

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Can negative eigenvalues be enforced such that $\Delta\tilde{H}(x) \leq -2\lambda$? **YES!**

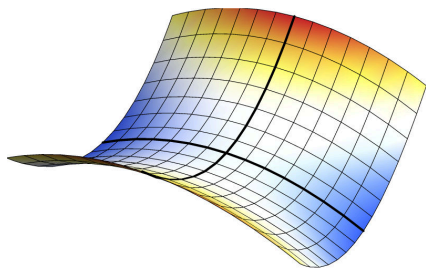


Figure : H around a saddle point

\tilde{H} is quadratic perturbation of H in $\sqrt{\varepsilon}$ -neighborhoods of critical points:

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Construction of Lyapunov function

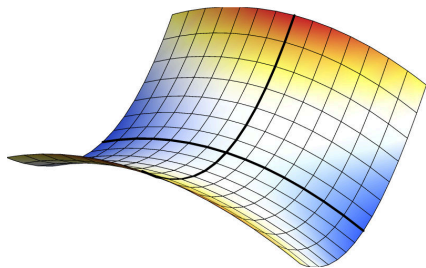


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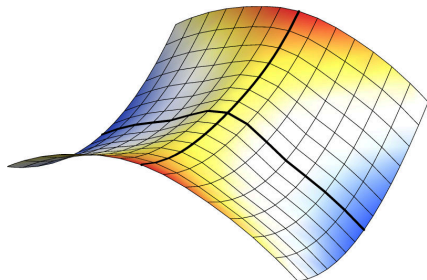


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Proof: Mean-difference estimate

Approximation step

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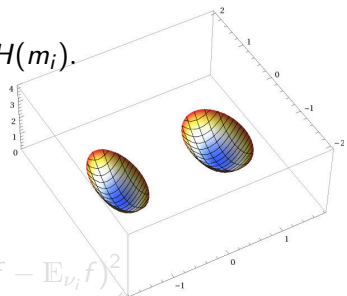
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Step 1: Approximate μ_0 and μ_1 by truncated Gaussians ν_0 and ν_1 :

$$\nu_i \sim \mathcal{N}(m_i, \varepsilon \Sigma_i) \llcorner B_{\sqrt{\varepsilon}}(m_i) \text{ with } \Sigma_i^{-1} := \nabla^2 H(m_i).$$

Introduce ν_0 and ν_1 as **coupling**:

$$\begin{aligned} (\mathbb{E}_{\mu_0} f - \mathbb{E}_{\mu_1} f)^2 &\leq (1 + \tau) \underbrace{(\mathbb{E}_{\nu_0} f - \mathbb{E}_{\nu_1} f)^2}_{\text{transport argument}} \\ &\quad + 2(1 + \tau^{-1}) \sum_{i=\{0,1\}} \underbrace{(\mathbb{E}_{\mu_i} f - \mathbb{E}_{\nu_i} f)^2}_{\text{approximation bound}} \end{aligned}$$



Approximation

bound follows from local PI and local LSI.

Approximation step

Goal: Find a good estimate for C in

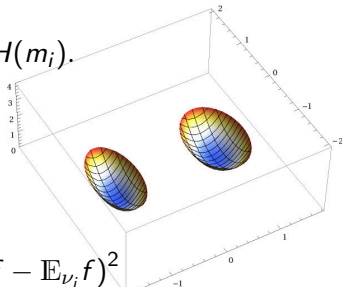
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Proof: Mean-difference estimate

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Sideremark: Weighted transport distance

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For $\nu_0, \nu_1 \ll \mu$ define the **weighted transport distance** by

$$\mathcal{T}_\mu^2(\nu_0, \nu_1) = \inf_{\{\Phi_s\}} \int \left| \int_0^1 \dot{\Phi}_s \circ \Phi_s^{-1} \frac{d\nu_s}{d\mu} ds \right|^2 d\mu.$$

$(\Phi_s)_{s \in [0,1]}$ is absolutely continuous in s : $(\Phi_s)_\# \nu_0 = \nu_s$.

Mean-difference revisited: Identify $\int |\nabla f|^2 d\mu = \|f\|_{\dot{H}^1(\mu)}^2$, then

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Indeed, it holds: $\mathcal{T}_\mu^2(\nu_0, \nu_1) = \|\nu_0 - \nu_1\|_{\dot{H}^{-1}(\mu)}^2$.

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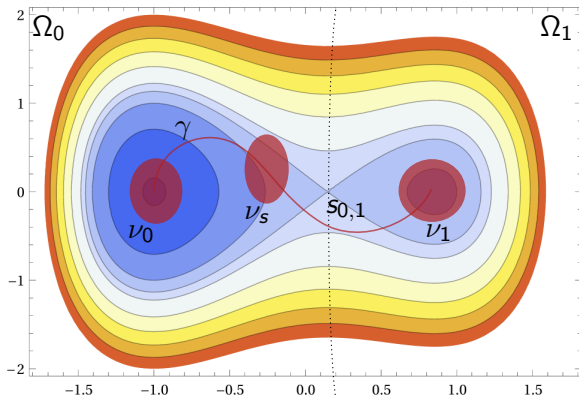
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Proof: Mean-difference estimate

Construction of transport interpolation

Step 3: Ansatz Φ_s such that $\nu_s = (\Phi_s)_\# \nu_0 = \mathcal{N}(\gamma_s, \Sigma_s) \llcorner B_{\sqrt{\varepsilon}}(\gamma_s)$

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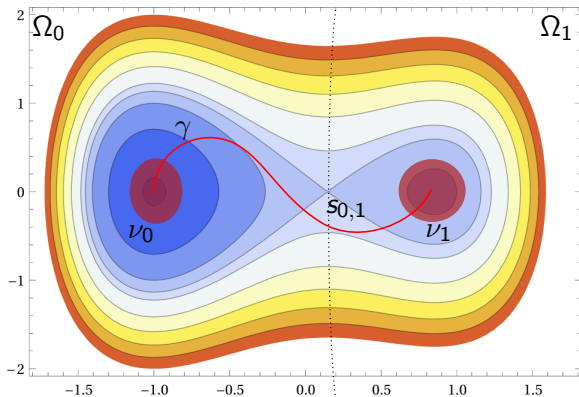


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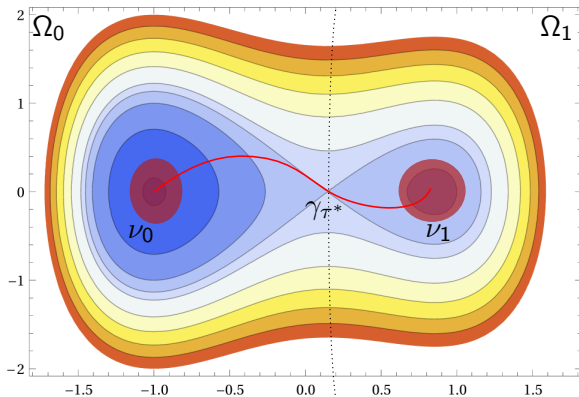
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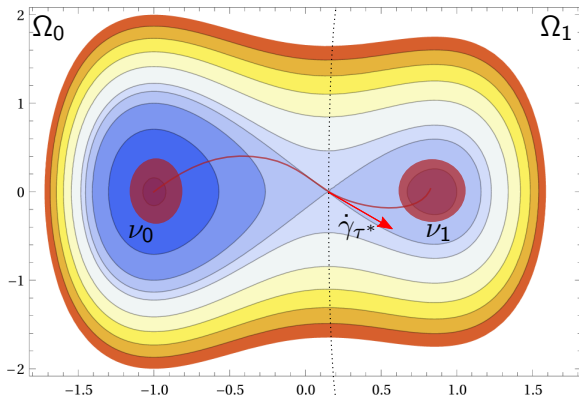
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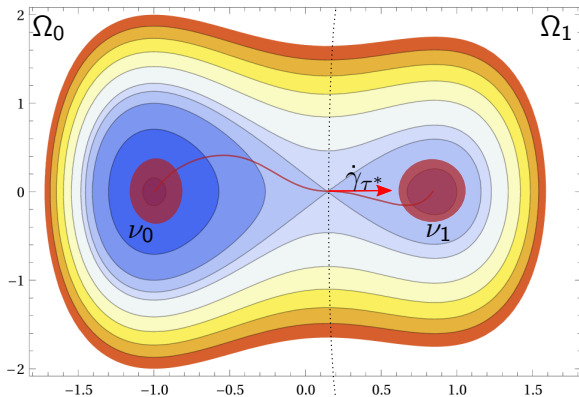


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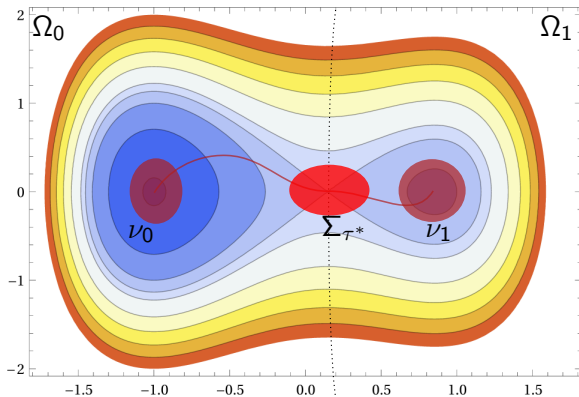
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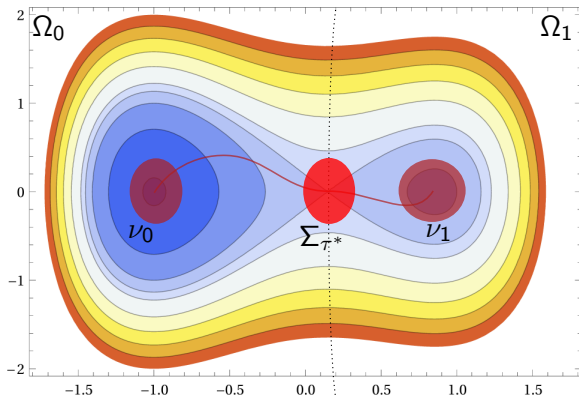
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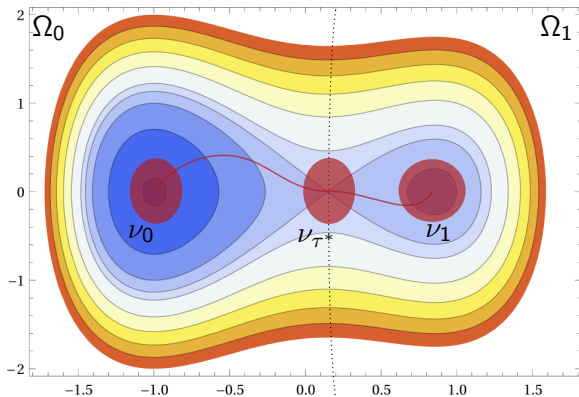
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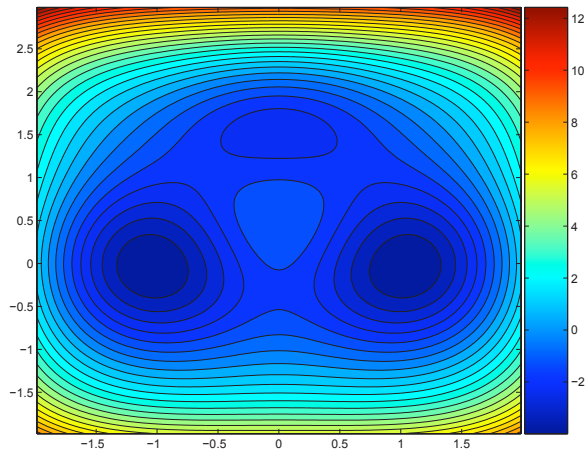
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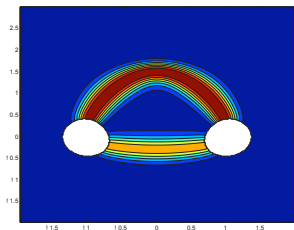
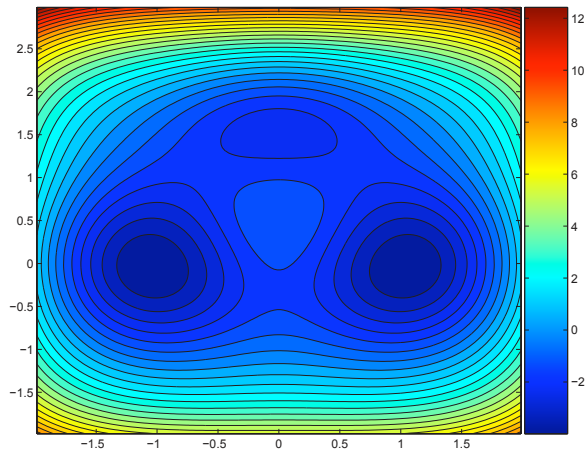


Observation of Eric Vanden-Eijden:
Compare flux between global minima at low
and very low temperature



Application to entropic switching

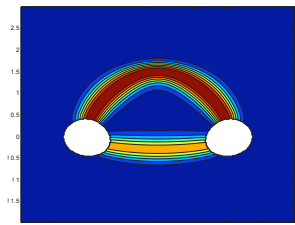
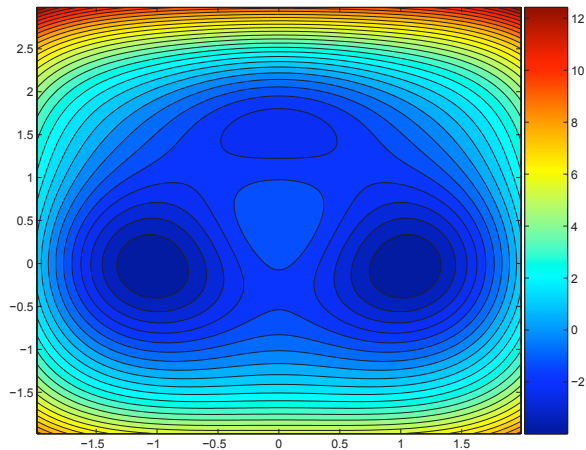
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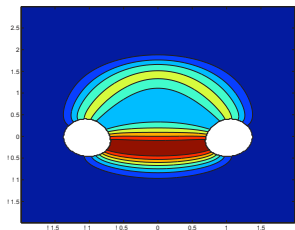
Very low temperature:
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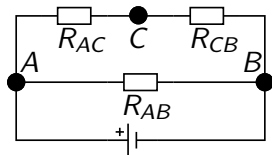


Very low temperature:
 $\varepsilon = 0.15$



Low temperature: $\varepsilon = 0.6$

- minima A, B, C become nodes
- saddle points become resistors
- reaction rate k_{AB} becomes the total conductance $R_{\text{tot}}^{-1}(A, B)$ between A, B .



This identification can be justified using the weighted transport distance

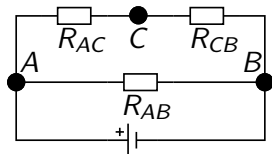
$$k_{AB} \approx \frac{1}{T_{\mu}^2(\mu_{\text{L}}A, \mu_{\text{L}}B)} \approx \frac{1}{R_{AB}} + \frac{1}{R_{AC} + R_{CB}},$$

where

$$R_{AB} = \inf_{\Phi \in \Pi(s_{AB})} \int \left| \int_0^1 \dot{\Phi}_s \circ \Phi_s^{-1} \frac{d\nu_s}{d\mu} ds \right|^2 d\mu.$$

$\Pi(s_{AB})$: transport interpolations between $\mu_{\text{L}}A$ and $\mu_{\text{L}}B$ across s_{AB} .

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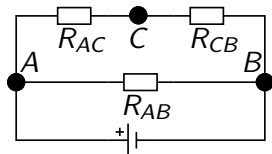
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	$\varepsilon = 0.15$	$\varepsilon = 0.6$
[MSVE06] TPT, flux	9.47×10^{-8}	1.912×10^{-2}
[MSVE06] TPT, committor	9.22×10^{-8}	1.924×10^{-2}
numerical simulation		$(1.918 \pm 0.052) \times 10^{-2}$
$\mathcal{T}(\nu_A, \nu_B), Z_\mu$ Gaussian	9.63×10^{-8}	2.373×10^{-2}
$\mathcal{T}(\nu_A, \nu_B), Z_\mu$ numerical	9.33×10^{-8}	1.926×10^{-2}

Table : Numerical estimates on $k_{A,B}$

[MSVE06] P. Metzner, C. Schütte, and E. Vanden-Eijnden, *Illustration of transition path theory on a collection of simple examples*. The Journal of chemical physics, 125:8, 2006.

- Functional inequalities quantify convergence to equilibrium
- Partitions and splitting induced from dynamic (two scales)
- Eyring-Kramers formula follows from two ingredients:
 - ▶ *good* local mixing
⇒ Lyapunov technique handles non-convex situations
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 - ▶ *sharp* estimates of mean-difference
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- Functional inequalities quantify convergence to equilibrium
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